Bounds on the Truncation Error by Finite Differences for the Goursat Problem

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1. Introduction. A common method for obtaining an approximate solution to the various boundary value problems for partial differential equations is by finite differences. The fundamental problem concerning such finite difference approximations is to show that the truncation error (i.e., the difference between the exact solution and the solution of the approximating finite difference problem) tends to zero with diminishing mesh size. Further the success of a finite difference method is related to the possibility of deducing explicit *a priori* bounds for the solution, or to the stability of the difference scheme. Various authors have studied the question of stability and *a priori* estimates connected with finite difference schemes concerning hyperbolic partial differential equations (see e.g. Richtmyer [14], with references).

In general the main tool used in obtaining estimates for the solution of hyperbolic equations is the energy method.

In this paper we obtain explicit bounds in terms of the data of the problem for the linear and non-linear second order partial differential equation of hyperbolic type in two independent variables, i.e. linear boundary value problem 2.1, 2.2 and non linear problem 4.1, 4.2. The present approach is based not on the energy method but on the method of majorants which is embodied in the statement of theorem 2.1 and its corollaries.

In section 2 we first define a suitable finite difference analogue of 2.1 and develop the finite difference analogue of Riemann's function and Riemann's formula for equation 2.1. Making use of this representation of the solution U(x, y) of boundary value problem 2.9, 2.10 and the majorizing theorem 2.1, together with its corollaries, we obtain in section 3 explicit bounds of $O(h^2)$ for the truncation errors |u(x, y) - U(x, y)|, $|u_x(x, y) - U(x, y)_x|$, $|u_y(x, y) - U(x, y)_y|$, $|u_{xy}(x, y) - U(x, y)_{xy}|$, where u and U denote the solution of 2.1, 2.2 and 2.9, 2.10 respectively. In section 4, employing the results and techniques of section 3, we obtain a similar bound of $O(h^2)$ for the truncation error E(x, y). for 4.1. Both bounds in section 3 and in section 4 are of nondecaying exponential nature and this exponential character makes our bounds rather poor when applied to problems whose solutions have decaying exponential behavior (i.e., $u_{xy} + u_x + u_y + u = 0$, u(x, 0) = x, u(0, y) = y, $u(x, y) = xe^{-y} + ye^{-x}$).

Applying Dames [5] criterion for stability, it is easily seen that the finite difference scheme considered is unconditionally stable.

In section 5 we give a simple numerical example where the bound for the truncation error is computed.

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2. Finite Difference Analogue. We shall pose a finite difference analogue of the characteristic boundary value problem given by

(2.1)
$$Lu \equiv u_{xy} + au_x + bu_y + cu = f \text{ in } R$$
$$u(x, 0) = \varphi(x)$$
$$u(0, y) = \psi(y)$$
$$\varphi(0) = \psi(0)$$

Where R is the first quadrant of the (x, y) plane. Let R_h be the set of grid points (mh, nh); m, n positive integers and let C_h be the boundary mesh points (mh, 0) and (0, nh). The real number n > 0 is called the "mesh constant". Any function W(x, y) which is defined at points of $R_h + C_h$ is called a "mesh function". We assume that the domain of definition of every mesh function is extended to the points $(mh - \frac{1}{2}h, nh), (mh, nh - \frac{1}{2}h), (mh - \frac{1}{2}h, nh - \frac{1}{2}h)$ by the following rules.

$$W(mh - \frac{1}{2}h, nh) = \frac{1}{2}[W(mh, nh) + W(mh - h, nh)]$$

$$W(mh, nh - \frac{1}{2}h) = \frac{1}{2}[W(mh, nh) + W(mh, nh - h)]$$

$$W(mh - \frac{1}{2}h, nh - \frac{1}{2}h) = \frac{1}{2}[W(mh - \frac{1}{2}h, nh) + W(mh - \frac{1}{2}h, nh - h)]$$

$$= \frac{1}{2}[W(mh, nh - \frac{1}{2}h) + W(mh - h, nh - \frac{1}{2}h)]$$

$$= \frac{1}{4}[W(mh, nh) + W(mh - h, nh) + W(mh - h, nh - h)]$$

By defining W(x, y) at such points as the averages of W(x, y) at nearby points of $R_h + C_h$ we make it possible to use the finite difference equivalent of the chain rule of differentiation. The desirability of this was pointed out by B. Wendroff [20]. At points of the extended domain of definition of W(x, y) we define the finite difference equivalent of partial differentiation by

(2.4)
$$W(x, y)_{x} = h^{-1}[W(x + \frac{1}{2}h, y) - W(x - \frac{1}{2}h, y)]$$
$$W(x, y)_{y} = h^{-1}[W(x, y + \frac{1}{2}h) - W(x, y - \frac{1}{2}h)].$$

From (2.4) we see that

$$W(x, y)_{XY} = W(x, y)_{YX} = h^{-1}[W(x, y + \frac{1}{2}h)_X - W(x, y - \frac{1}{2}h)_X]$$

$$= h^{-1}[W(x + \frac{1}{2}h, y)_Y - W(x - \frac{1}{2}h, y)_Y]$$

$$= h^{-2}[W(x + \frac{1}{2}h, y + \frac{1}{2}h) - W(x + \frac{1}{2}h, y - \frac{1}{2}h)$$

$$- W(x - \frac{1}{2}h, y + \frac{1}{2}h) + W(x - \frac{1}{2}h, y - \frac{1}{2}h)].$$

We note that for the mesh functions V(x, y), W(x, y) and m, n integers we have

$$[V(mh - \frac{1}{2}h, y)W(mh - \frac{1}{2}h, y)]_{x}$$

$$= V(mh - \frac{1}{2}h, y)W(mh - \frac{1}{2}h, y)_{x} + V(mh - \frac{1}{2}h, y)_{x}W(mh - \frac{1}{2}h, y)$$

$$[V(x, nh - \frac{1}{2}h)W(x, nh - \frac{1}{2}h)]_{r}$$

$$= V(x, nh - \frac{1}{2}h)W(x, nh - \frac{1}{2}h)_{r} + V(x, nh - \frac{1}{2}h)_{r}W(x, nh - \frac{1}{2}h)$$

Let the mesh functions A and B be given at points $(x + \frac{1}{2}h, y)$ and $(x, y + \frac{1}{2}h)$ respectively by

(2.7)
$$A(x + \frac{1}{2}h, y) = a(x + \frac{1}{2}h, y)$$
$$B(x, y + \frac{1}{2}h) = b(x, y + \frac{1}{2}h)$$

where $x = mh - \frac{1}{2}h$, $y = nh - \frac{1}{2}h$. Similarly let Φ and Ψ be defined on their respective subsets of C_h as

(2.8)
$$\begin{aligned} \Phi(x) &= \varphi(x) \qquad x = mh \\ \Psi(y) &= \psi(y) \qquad y = nh. \end{aligned}$$

We formulate the following finite difference analogue of the characteristic boundary value problem (2.1) and (2.2).

(2.9)

$$L_h U \equiv U_{xy} + A U_x + B U_y + C U = f(x, y)$$

$$x = mh - \frac{1}{2}h, \quad y = nh - \frac{1}{2}h; \quad m, n \text{ positive integers},$$

$$U(mh, 0) = \Phi(mh)$$
(2.10)

$$U(0, nh) = \Psi(nh).$$

In (2.9) A and B are given as averages of the values at the points of (2.7) whereas C(x, y) and f(x, y) are values at the point (x, y). We see that (2.9) and (2.10) is a system of linear equations for U(mh, nh) and that in each equation of (2.9) the mesh function U(mh, nh) is given as a linear combination of U(mh - h, nh - h), U(mh, nh - h), U(mh - h, nh) and F. Hence (2.9) and (2.10) can be solved explicitly.

In preparation for deriving a finite difference form of Riemann's formula we now introduce a certain Riemann-Stieltjes integral as has been done in elliptic problems by G. Forsythe and W. Wasow [8]. Let the step functions $\mu(x)$, $\nu(y)$ be defined by

(2.11)
$$\mu(x) = mh; \qquad mh - \frac{1}{2}h < x \le mh + \frac{1}{2}h \\ \nu(y) = nh; \qquad nh - \frac{1}{2}h < y \le nh + \frac{1}{2}h.$$

The Riemann-Stieltjes integral with respect to $\mu(x)$ is seen to be

(2.12)
$$\int_0^{mh} f(x) \ d\mu = h \sum_{m=1}^M f(mh - \frac{1}{2}h).$$

In particular if W(x, y) is a mesh function and we assume that the definition of W(x, y) is extended in a smooth manner to the entire (x, y) plane then we have

(2.13)
$$\int_0^{Mh} W(x,y) \ d\mu = h \sum_{m=1}^M W(mh - \frac{1}{2}h, y).$$

Likewise the Riemann-Stieltjes integral with respect to $\nu(y)$ is seen to be

(2.14)
$$\int_0^{Nh} W(x, y) \, d\nu = h \sum_{n=1}^N W(x, nh - \frac{1}{2}h).$$

Let $\xi = Mh$. $\eta = Nh$; M, N positive integers and assume that W(x, y) and

V(x, y) are mesh functions such that W vanishes at points of the x and y axes. From the definition of W(x, y), V(x, y) in the extended domain together with the chain rule for differences (2.6) we can carry through the usual integration by parts with differentiation replaced by differencing to obtain

(2.15)

$$\int_{0}^{\xi} d\mu \int_{0}^{\eta} d\nu V(x, y) L_{h} W(x, y) = \int_{0}^{\xi} d\mu \int_{0}^{\eta} d\nu W(x, y) L_{h}^{*} V(x, y)$$

$$- \int_{0}^{\xi} [V(x, \eta)_{x} - B(x, \eta) V(x, \eta)] W(x, \eta) d\mu$$

$$- \int_{0}^{\eta} [V(\xi, y)_{y} - A(\xi, y) V(\xi, y)] W(\xi, y) d\nu + V(\xi, \eta) W(\xi, \eta).$$

In (2.15) we have introduced the finite difference adjoint operator L_h^* of L_h which is given by

(2.16)
$$L_h^* V(x, y) \equiv V_{XY} - (AV)_X - (BV)_Y + CV.$$

We now define for each $P = (\xi, \eta)$ the mesh function $V(x, y; P) = V(x, y; \xi, \eta)$ to be the solution of the finite difference problem

$$L_{h}^{*}V(x, y; P) = 0; \quad 0 < x = mh - \frac{1}{2}h < \xi \text{ and}$$

$$0 < y = nh - \frac{1}{2}h < \eta$$

$$(2.17) \quad V(x, \eta; P)_{x} - B(x, \eta)V(x, \eta; P) = 0; \quad 0 < x = mh - \frac{1}{2}h < \xi$$

$$V(\xi, y; P)_{x} - A(\xi, y)V(\xi, y; P) = 0; \quad 0 < y = nh - \frac{1}{2}h < \eta$$

$$V(P; P) = 1.$$

Again we observe that V(x, y; P) is given explicitly by (2.17) so that there is no question as to its existence and uniqueness. Also for V determined by (2.17) when substituted into (2.15) we have the finite difference form of Riemann's formula

(2.18)
$$W(\xi,\eta) = \int_0^{\xi} d\mu \int_0^{\eta} d\nu V(x,y;\xi,\eta) L_h W(x,y)$$

for mesh functions $W(\xi, \eta)$ which vanish at points of the x-axis and the y-axis. If W(P) does not vanish on the axes then (2.18) takes the more general form

(2.19)

$$W(\xi,\eta) = +V(0,0;\xi,\eta)W(0,0) + \int_{0}^{\xi} [W(x,0)_{x} - B(x,0)W(x,0)]V(x,0) d\mu + \int_{0}^{\eta} [W(0,y)_{x} - A(0,y)W(0,y)]V(0,y) d\nu + \int_{0}^{\xi} d\mu \int_{0}^{\eta} d\nu V(x,y;\xi,\eta)L_{h} W(x,y).$$

We shall now derive some further finite difference analogues of classical results concerning the solution of (2.1), (2.2). Finally we shall conclude this section with a theorem concerning certain majorizing finite difference problems. We first give a

finite difference form of the function $e^{f(x)}$. This function arises as a solution of the differential equation

(2.20)
$$y'(x) = f'(x)y(x).$$

The finite difference exponential arises as the solution of the finite difference analogue of (2.20) as given in the following lemma.

LEMMA 2.1. For K(x) given at x = mh; $m = 0, 1, \dots, M$ we define the mesh function $e_h^{K(x)}$ as the solution of the difference equation

(2.21)
$$[e_h^{K(x)}]_x = K(x)_x e_h^{K(x)}; \quad x = \frac{1}{2}h, \frac{3}{2}h, \cdots, (M - \frac{1}{2})h.$$

In (2.21) we define

$$K(mh - \frac{1}{2}h) = \frac{1}{2}[K(mh) + K(mh - h)]$$
$$e_{h}^{K(mh - \frac{1}{2}h)} = \frac{1}{2}[e_{h}^{K(mh)} + e_{h}^{K(mh - h)}].$$

Then for x = mh, $e_h^{K(x)}$ is given by the formula

(2.22)
$$e_{h}^{K(x)} = e_{h}^{K(0)} \prod_{m=1}^{M} \left[\frac{1 + \frac{1}{2}hK(mh - \frac{1}{2}h)_{x}}{1 - \frac{1}{2}hK(mh - \frac{1}{2}h)_{x}} \right]$$

Proof. For fixed m we solve (2.21) to obtain

$$e_{h}^{K(mh)} = e_{h}^{K(mh-h)} \left[\frac{1 + \frac{1}{2}hK(mh - \frac{1}{2}h)_{X}}{1 - \frac{1}{2}hK(mh - \frac{1}{2}h)_{X}} \right]$$

Solving this recursion formula we obtain (2.22).

We can apply Lemma 2.1 to the boundary conditions in (2.17) to obtain

$$V(\bar{x},\eta;\xi,\eta) = \exp_{h}\left[-\int_{\bar{x}}^{\xi} B(x,\eta)d\mu\right] \qquad \qquad \bar{x} \leq \xi$$

(2.23)
$$V(\xi, \bar{y}; \xi, \eta) = \exp_{h}\left[-\int_{\bar{y}}^{\eta} A(\xi, y) d\nu\right] \qquad \bar{y} \leq \eta$$

since $V(\xi, \eta; \xi, \eta) = 1$.

A comparison of the series expansion of e^x and e_h^x reveals that

$$(2.24) e^x \leq e_h^x x = mh \geq 0.$$

We have the following inequality relating our Stieltjes integral to the Riemann integral.

LEMMA 2.2. If $f(x) \in C^2$ and $f''(x) \ge 0$, then for any positive integer m we have

(2.25)
$$\int_{0}^{mh} f(x) \ d\mu \leq \int_{0}^{mh} f(x) \ dx.$$

Proof. If $x_0 = rh - \frac{1}{2}h$, r an integer, then

(2.26)
$$f(x) = f(x_0) + f'(x_0)(x - x_0) + \int_{x_0}^{x} f''(\xi)(x - \xi) d\xi.$$

Hence

(2.27)
$$\int_{rh-h}^{rh} f(x) \, dx = hf(x_0) + \int_{rh-h}^{rh} dx \int_{x_0}^{x} f''(\xi)(x-\xi) \, d\xi.$$

Since the second term on the right hand side of (2.27) is positive we see that

(2.28)
$$\int_{\tau h-h}^{\tau h} f(x) \ dx \ge h f(x_0) = \int_{\tau h-h}^{\tau h} f(x) \ d\mu$$

Summing the inequality (2.28) over integral values of r yields (2.25).

We have the following lemma regarding differencing with respect to a parameter. LEMMA 2.3. If W(x, y) is a mesh function then for $y = mh - \frac{1}{2}h$

$$\left[\int_0^y W(x,y) \ d\mu\right]_Y = W(y,y+\frac{1}{2}h) + \int_0^{y-\frac{1}{2}h} W(x,y)_Y \ d\mu$$

Proof.

$$\begin{bmatrix} \int_0^y W(x, y) \ d\mu \end{bmatrix}_r = h^{-1} \int_0^{y+\frac{1}{2}h} W(x, y + \frac{1}{2}h) \ d\mu - \int_0^{y-\frac{1}{2}h} W(x, y - \frac{1}{2}h) \ d\mu$$
$$= h^{-1} \int_{y-\frac{1}{2}h}^{y+\frac{1}{2}h} W(x, y + \frac{1}{2}h) \ d\mu + \int_0^{y-\frac{1}{2}h} W(x, y)_r \ d\mu = W(y - \frac{1}{2}h, y + \frac{1}{2}h)$$
$$+ \int_0^{y-\frac{1}{2}h} W(x, y)_r \ d\mu$$

Since we shall be using the method of majorants to bound the "discretization error" (the difference between the solutions of (2.1) and (2.9)) the following theorem on majorizing problems for (2.9) is of central importance.

THEOREM 2.1. Let W(x, y) be a solution of the finite difference problem

(2.29)
$$W_{XY} = C_1 W_X + C_2 W_Y + C_3 W + C_4$$

at points $(mh - \frac{1}{2}h, nh - \frac{1}{2}h)$ where the functions $C_i(x, y) \ge 0$. We further assume that

(2.30)
$$\begin{array}{ll} W(x,0) \geq 0, & W(x,0)_x \geq 0, & x \geq 0 \\ W(0,y) \geq 0, & W(0,y)_y \geq 0, & y \geq 0. \end{array}$$

Finally let

$$(2.31) (1-\delta_3) > 0, (1-\delta_4) \ge 0$$

where

$$\delta_{1} = \frac{1}{2}h\left(C_{1} + \frac{3C_{3}h}{4}\right)$$

$$\delta_{2} = \frac{1}{2}h\left(C_{2} + \frac{3C_{3}h}{4}\right)$$

$$\delta_{3} = \frac{1}{2}h\left(C_{1} + \frac{C_{3}h}{4}\right)$$

$$\delta_{4} = \frac{1}{2}h\left(C_{2} + \frac{C_{3}h}{4}\right)$$

$$\delta_{5} = h[C_{3}W(x, y) + C_{4}]$$

Then we conclude that

$$(2.33) W \ge 0, W_x \ge 0, W_y \ge 0, W_{xy} \ge 0$$

at points of R_n .

Proof. The theorem is proved by induction. Let x = mh, y = nh be an arbitrary point of R_h . We introduce the notation

$$a_1 = W(x, y)_x$$
, $a_2 = W(x, y + h)_x$,
 $b_1 = W(x, y)_r$, $b_2 = W(x + h, y)_r$, $d_1 = W(x, y)$

Our induction assumption is that

$$(2.34) a_1 \ge 0, b_1 \ge 0, d_1 \ge 0$$

We see that 2.29 can be written in the alternate forms

Solving these equations simultaneously under the assumption (2.31) we see that

$$a_{2}\left(1-\delta_{3}-\frac{\delta_{1}\delta_{4}}{1-\delta_{4}}\right) = a_{1}\left(1+\delta_{1}+\frac{\delta_{1}\delta_{4}}{1-\delta_{4}}\right)$$
$$+ b_{1}\left[\delta_{2}+\delta_{4}\left(\frac{1+\delta_{2}}{1-\delta_{4}}\right)\right] + \delta_{5}\left(1+\frac{\delta_{4}}{1-\delta_{4}}\right)$$
$$b_{2}\left(1-\delta_{4}-\frac{\delta_{2}\delta_{3}}{1-\delta_{3}}\right) = b_{1}\left(1+\delta_{2}+\frac{\delta_{2}\delta_{3}}{1-\delta_{3}}\right)$$
$$+ a_{1}\left[\delta_{1}+\delta_{3}\left(\frac{1+\delta_{1}}{1-\delta_{3}}\right)\right] + \delta_{5}\left(1+\frac{\delta_{3}}{1-\delta_{3}}\right).$$

From (2.35) and (2.34) it follows immediately that

$$a_2 \ge a_1 \ge 0; \qquad b_2 \ge b_1 \ge 0$$

and hence that

$$W(x+h, y+h) \ge \begin{cases} W(x+h, y) \ge 0\\ \\ W(x, y+h) \ge 0 \end{cases}$$

Clearly then $W(x + \frac{1}{2}h, y + \frac{1}{2}h)_{xy} \ge 0$ and the theorem is proved.

We can obtain an explicit upper bound for W(x, y) in the following special case.

COROLLARY 1. Let W(x, y) satisfy the hypotheses of Theorem 2.1. In addition we assume that C_i are constants and that W(x, 0) = W(0, y) = 0. Then W(x, y) satisfies the inequality

(2.36)
$$W(\xi,\eta) \leq C_4 \int_0^{\xi} dx \int_0^{\eta} dy \left[e_h \sqrt{C_5} \{ (\xi-x) + (\eta-y) \} \right] \left[e_h^{C_2(\xi-x) + C_1(\eta-y)} \right]$$

where

$$(2.37) C_5 = C_3 + C_1 C_2 \,.$$

Proof. If we define the function V(x, y) at points x = mh, y = nh by the equation

(2.38)
$$W(x, y) = V(x, y)e_h^{c_2 x + c_1 y}$$

we see that V(x, y) is a solution of the finite difference problem

(2.39)
$$V_{xy}(x,y) = C_5 V(x,y) + C_4 e_h^{-(C_2 x + C_1 y)}, \quad x = mh - \frac{1}{2}h, y = nh - \frac{1}{2}h$$
$$V(x,0) = V(0,y) = 0$$

The finite difference Riemann's function, $V^*(x, y; \xi, \eta)$, related to the problem (2.39) as given in (2.17) satisfies the adjoint problem

(2.40)
$$V^{*}(x, y; \xi, \eta)_{XY} = C_{5}V^{*}(x, y; \xi, \eta)$$
$$V^{*}(x, \eta; \xi, \eta)_{X} = 0$$
$$V^{*}(\xi, y; \xi, \eta)_{Y} = 0$$

It is easily verified that

(2.4)
$$V^*(x, y; \xi, \eta) \leq e_h \sqrt{C_*} [(\xi - x) + (\eta - y)]$$

From the finite difference Riemann's formula (2.18) with L_h given in (2.39) we see that

(2.42)
$$V(\xi,\eta) \leq C_4 \int_0^{\xi} d\mu \int_0^{\eta} d\nu \left[e_h \sqrt{c_5} \{ (\xi-x) + (\eta-y) \} \right] \left[e_h^{-(C_2 x + C_1 y)} \right]$$

The result (2.36) now follows from (2.38) and Lemma 2.2. The substitution (2.38) and the subsequent bound are motivated by analogous results for the differential equation (2.1), R. von Mises [17, p. 812].

The following corollary gives the form of the majorizing theorem which we shall utilize in obtaining bounds on the discretization error.

COROLLARY 2. Let the quantities in Theorem 2.1 have superscripts (1) and (2) to indicate two distinct finite difference problems. We replace the assumption $C_i^{(1)}(x, y) \ge 0$ by $C_i^{(2)} \ge |C_i^{(1)}|$ and (2.30) by $W^{(i)}(x, 0) = W^{(i)}(0, y) = 0$. Then it follows that

(2.43)
$$\begin{array}{ll} W^{(2)}(x,y) & \geq \mid W^{(1)}(x,y) \mid; & W^{(2)}(x,y)_{\mathbf{X}} \geq \mid W^{(1)}(x,y)_{\mathbf{X}} \mid \\ & W^{(2)}(x,y)_{\mathbf{Y}} \geq \mid W^{(1)}(x,y)_{\mathbf{Y}} \mid; & W^{(2)}(x,y)_{\mathbf{X}\mathbf{Y}} \geq \mid W^{(1)}(x,y)_{\mathbf{X}\mathbf{Y}} \mid \end{array}$$

Proof. As in the proof of Theorem 2.1 the proof is by induction. From the hypotheses of the corollary together with (2.35) and the induction assumption we see that

$$\begin{aligned} |a_{2}^{(1)}| &\leq \left[1 - |\delta_{3}^{(1)}| - \frac{|\delta_{1}^{(1)}\delta_{4}^{(1)}|}{1 - |\delta_{4}^{(1)}|}\right]^{-1} \left\{ |a_{1}^{(1)}| \left(1 + |\delta_{1}^{(1)}| \right) \\ (2.44) &+ \frac{|\delta_{1}^{(1)}\delta_{4}^{(1)}|}{1 - |\delta_{4}^{(1)}|} \right) + |b_{1}^{(1)}| \left[|\delta_{2}^{(1)}| + |\delta_{4}^{(1)}| \frac{1 + |\delta_{2}^{(1)}|}{1 - |\delta_{4}^{(1)}|} \right] \\ &+ |\delta_{5}^{(1)}| \left(1 + \frac{|\delta_{4}^{(1)}|}{1 - |\delta_{4}^{(1)}|} \right) \right] \leq a_{2}^{(2)}.\end{aligned}$$
Similarly

Similarly

$$(2.45) | b_2^{(1)} | \le b_2^{(2)}$$

From (2.44) and (2.45) it is clear that

$$(2.46) | W^{(1)}(x+h, y+h) | \leq | W^{(1)}(x, y) | + | a_1^{(1)} | + | b_{(2)}^{(1)} | \\ \leq W^{(2)}(x+h, y+h)$$

and finally $|W_{XY}^{(1)}| \leq W_{XY}^{(2)}$ follows from the difference equation (2.29).

3. Bounds for the Discretization Error. In Theorem 3.1 we shall obtain an explicit bound on the discretization error in terms of the data and certain higher derivatives of the solution. Here corollaries 1 and 2 of Theorem 2.1 form the basis for the error estimate. Following this we show how to bound the higher derivatives of u(x, y) in terms of data.

THEOREM 3.1. Let u(x, y) be the solution of (2.1) and (2.2) with $\varphi(x) = \psi(y) = 0$ and let U(x, y) be the solution of (2.9) and (2.10). Define

(3.1)

$$K_{1}(\xi,\eta) = \max_{x \in [0,\xi], y \in [0,\eta]} |a(x,y)|$$

$$K_{2}(\xi,\eta) = \max_{x \in [0,\xi], y \in [0,\eta]} |b(x,y)|$$

$$K_{3}(\xi,\eta) = \max_{x \in [0,\xi], y \in [0,\eta]} |c(x,y)|$$

$$K_{4}(\xi,\eta) \ge \max_{x \in [0,\xi], y \in [0,\eta]} |L_{h} u(x,y)|.$$

Then the discretization error E(x, y) = u(x, y) - U(x, y) has the bound

$$(3.2) |E(\xi,\eta)| \leq K_4(\xi,\eta) \int_0^{\xi} dx \int_0^{\eta} dy \, [e_h \sqrt{c_5} [(\xi-x)+(\eta-y)]] [e_h^{K_2(\xi,\eta)}(\xi-x)+K_1(\xi,\eta)(y-\eta)]$$

where $K_4(\xi, \eta)$ is given by (3.5). For $0 \leq \bar{x} \leq \xi$, $0 \leq \bar{y} \leq \eta$ the bound for $|E(\bar{x}, \bar{y})|$ is given by (3.2).

Proof.

By Corollary 1 and Corollary 2 of Theorem 2.1 we see that (3.2) is valid for any K_4 which satisfies (3.1).

Before obtaining the bound (3.5) for $K_4(\xi, \eta)$ we introduce the notation

(3.3)
$$\varphi_{M} = \max_{x \in [0,\xi], y \in [0,\eta]} \varphi(x,y) \qquad \varphi_{m} = \min_{x \in [0,\xi], y \in [0,\eta]} \varphi(x,y)$$

for an arbitrary function (x, y). Let $x = mh - \frac{1}{2}h$, $y = nh - \frac{1}{2}h$, then by Taylor's theorem

$$\begin{aligned} | u(x, y)_{XY} - u(x, y)_{xy} | &\leq (\frac{1}{2}h)^2 \frac{1}{4!} \{ | u_{xxxx} |_M + 4 | u_{xxxx} |_M + 6 | u_{xxyy} |_M \\ &+ 4 | u_{xyyy} |_M + | u_{yyyyy} |_M \} = \alpha_1 h^2, \\ | u(x, y)_X - u(x, y)_x | &\leq \frac{1}{4}h^2 \{ | u_{xxx} |_M + 3 | u_{xxy} |_M + 3 | u_{xyy} |_M \\ (3.4) &+ | u_{yyy} |_M \} = h^2 \alpha_2, \\ | u(x, y)_Y - u(x, y)_y | &\leq h^2 \alpha_2, \\ | \frac{1}{4} \{ u(x - \frac{1}{2}h, y - \frac{1}{2}h) + u(x - \frac{1}{2}h, y + \frac{1}{2}h) \\ &+ a(x + \frac{1}{2}h, y - \frac{1}{2}h) + u(x + \frac{1}{2}h, y + \frac{1}{2}h) \} - u(x, y) \leq \frac{1}{8}h^2 \{ | u_{xx} |_M \\ &+ 2 | u_{xy} |_M + | u_{yy} |_M \} = h^2 \alpha_3. \end{aligned}$$

Certain of the derivatives on the right side of (3.4) can be replaced by $O(h^2)$ terms which involve even higher derivatives of U(x, y). It is now clear that

(3.5)
$$|L_h u - Lu| \leq h^2 \{\alpha_1 + (|a|_M + |b|_M)\alpha_2 + |c|_M \alpha_3 + \frac{1}{8} (|a_{xx}|_M |u_x|_M + |b_{yy}|_M |u_y|_M) \}.$$

In view of the definition of α , all that remains is to obtain bounds for the higher derivatives of u(x, y). These will now be bounded starting with the lowest order.

From Frank & v. Mises [17; p. 812] and the method of majorants we see that the Riemann's function of the majorizing problem is given by

(3.6)
$$V(x, y; \xi, \eta) = I_0 (2\sqrt{\beta(\xi - x)(\eta - y)}) e^{|b|_M (\xi - x) + |a|_M (\eta - y)}$$
$$\beta = |a|_M |b|_M + |c|_M.$$

Hence we have the inequalities

$$| u |_{\mathbf{M}} \leq \int_{0}^{\xi} dx \int_{0}^{\eta} dy V(x, y; \xi, \eta) | f(x, y) |$$

$$| u_{x} |_{\mathbf{M}} \leq \int_{0}^{\eta} dy V(\xi, y; \xi, \eta) | f(\xi, y) | + \int_{0}^{\xi} dx \int_{0}^{\eta} dy V_{\xi}(x, y; \xi, \eta) | f(x, y) |$$

$$| u_{y} |_{\mathbf{M}} \leq \int_{0}^{\xi} dx V(x, \eta; \xi, \eta) | f(x, \eta) | + \int_{0}^{\xi} dx \int_{0}^{\eta} dy V_{\eta}(x, y; \xi, \eta) | f(x, y) |$$

$$| u_{xy} |_{\mathbf{M}} \leq | a |_{\mathbf{M}} | u_{x} |_{\mathbf{M}} + | b |_{\mathbf{M}} | u_{y} |_{\mathbf{M}} + | c |_{\mathbf{M}} | u |_{\mathbf{M}} + | f |_{\mathbf{M}}.$$

The functions u_{xx} and u_{yy} are bounded as follows. We write the differential equation 2.1 in the form

(3.8)
$$u_{xy} + au_x = f - bu_y - cu \equiv g.$$
$$u_{xy} + bu_y = f - au_x - cu \equiv l.$$

The left side of (3.8, a) can be integrated with respect to y by using the integrating factor

(3.9)
$$e^{p(\xi,\eta,y)} \equiv \exp\left[\int_{y}^{\eta} a(\xi,\lambda) \ d\lambda\right].$$

This together with the assumption that $u_x(x, 0) = 0$ yields

(3.10)
$$u_{\xi}(\xi,\eta) = \int_{0}^{\eta} g(\xi,y) e^{p(\xi,\eta,y)} \, dy.$$

We have similar integral representation for $u_{\eta}(\xi, \eta)$. Differentiation of (3.10) yields

(3.11)
$$u_{\xi\xi}(\xi,\eta) = \int_0^{\eta} \{g(\xi,y)p_{\xi}(\xi,\eta,y) + g_{\xi}(\xi,y)\}e^p \, dy.$$

By repeating this process for successively higher derivatives and making use of the integral representations for lower order derivatives, we can express the derivatives of all orders in terms of the data. We can then apply well known inequalities to these representations to obtain computable bounds for α_1 , α_2 , α_3 , in (3.4).

We conclude this section with explicit bounds on the error between the first and

the second order mixed derivatives of u(x, y) and their finite difference counterparts of U(x, y).

We see that:

$$| u_{x}(P) - U_{x}(P) | \leq |E(P)_{x}| + |u_{x}(P) - u(P)_{x}|$$
(3.12)
$$| u_{y}(P) - U(P)_{y} | \leq |E(P)_{y}| + |u_{y}(P) - u(P)_{y}|$$

$$| u_{xy}(P) - U(P)_{xy} | \leq |E(P)_{xy}| + |u_{xy}(P) - u(P)_{xy}|.$$

The $O(h^2)$ bounds for $|u_x(P) - u(P)_x|$, $|u_y(P) - u(P)_y|$, and $|u_{xy}(P) - u(P)_{xy}|$, in terms of higher derivatives of u(x, y) are given by $h^2\alpha_2$, $h^2\alpha_2$ and $h^2\alpha_1$ respectively, where α_2 and α_1 are defined by (3.4).

To bound $|E(P)_{\mathbf{x}}|, |E(P)_{\mathbf{y}}|$ and $|E(P)_{\mathbf{xy}}|$, we consider the boundary value problem

(3.13)
$$E_{xy} = -AE_x - BE_y - CE + \sigma(x, y)$$
$$E(x, 0) = E(0, y) = 0$$

where $\sigma(x, y) = L^{h}u - Lu$ and $|\sigma(x, y)| \leq K_{5}$, where K_{5} denotes the right-hand side of (3.5).

From Corollary 2 of section 2 it follows that:

(3.14)
$$|E(x, y)| \leq E^*(x, y), \qquad |E(x, y)_x| \leq E^*(x, y)_x, \\ |E(x, y)_r| \leq E^*(x, y)_r, \qquad |E(x, y)_{xr}| \leq E^*(x, y)_{xr}$$

where $E^*(x, y)$ is the solution of the boundary value problem.

(3.15)
$$E^*(x, y)_{XY} = C_1 E_X^*(x, y) + C_2 E^*(x, y)_Y + C_3 E^*(x, y) + K_5, \\E^*(x, 0) = E^*(0, y) = 0$$

with $C_1 = \max_{x,y\in R} |A(x, y)|$, $C_2 = \max_{x,y\in R} |B(x, y)|$, and $C_3 = \max_{x,y\in R} |C(x, y)|$. By the transformation

(3.16)
$$E^*(x, y) = \bar{E}(x, y)e_{\hbar}^{C_2 x + C_1 y},$$

(3.15) reduces to

(3.17)
$$\bar{E}_{XY}(x, y) = C_5 \bar{E}(x, y) + K_5 e_h^{-(C_2 x + C_1 y)} \\ \bar{E}(x, 0) = \bar{E}(0, y) = 0,$$

with $C_5 = C_1 C_2 + C_3$.

Denoting by $\overline{V}(\overline{x}, \overline{y}, x, y)$ the Riemann function associated with (3.17) and making use of the symmetry property of the Riemann function [2, p. 454] we have

(3.18)
$$V(\bar{x}, \bar{y}; x, y)_{XY} = C_5 V(\bar{x}, \bar{y}; x, y)$$

$$V(\bar{x}, \bar{y}; \bar{x}, y) = V(\bar{x}, \bar{y}; x, \bar{y}) = 1$$

From (2.41) we have

(3.19)
$$\tilde{V}(\bar{x},\bar{y};x,y) \leq e_{h}^{\sqrt{C_{\mathfrak{s}}}[(x-\bar{x})+(y-\bar{y})]}.$$

Now let $z(\bar{x}, \bar{y}; x, y) = e_h^{\sqrt{C_b}[(x-\bar{x})+(y-\bar{y})]}$, and $\bar{z} = z - \bar{V}(\bar{x}, \bar{y}, x, y)$. It is easily

seen that

(3.20)
$$\bar{z}(\bar{x},\bar{y};x,y)_{xy} = C_5 \bar{z}(\bar{x},\bar{y};x,y) \\ \bar{z}(\bar{x},\bar{y};\bar{x},y) = e_h^{\sqrt{C_5}(y-\bar{y})} - 1, \bar{z}(\bar{x},\bar{y},x,\bar{y}) = e_h^{\sqrt{C_5}(x-\bar{x})} - 1.$$

From (3.20) and Theorem 2.1, it follows that

 $\bar{z}, \bar{z}_{x}, \bar{z}_{y}, \bar{z}_{xy} \geq 0,$

and hence

(3.21)

$$\bar{V}(\bar{x}, \bar{y}; x, y) \leq e_{h}^{\sqrt{C_{5}}(x-\bar{x}+y-\bar{y})}$$

$$\bar{V}(\bar{x}, \bar{y}; x, y)_{\mathbf{x}} \leq \sqrt{C_{5}} e_{h}^{\sqrt{C_{5}}(x-\bar{x}+y-\bar{y})}$$

$$\bar{V}(\bar{x}, \bar{y}; x, y)_{\mathbf{x}} \leq \sqrt{C_{5}} e_{h}^{\sqrt{C_{5}}(x-\bar{x}+y-\bar{y})}$$

$$V(\bar{x}, \bar{y}; x, y)_{\mathbf{x}\mathbf{y}} \leq C_{5} e_{h}^{\sqrt{C_{5}}(x-\bar{x}+y-\bar{y})}$$

Now by (2.18), $\overline{E}(x, y)$, $\overline{E}(x, y)_x$, $\overline{E}(x, y)_y$, and $\overline{E}(x, y)_{xy}$ may be written as

$$\begin{split} \bar{E}(x,y) &= K_5 \int_0^x d\mu \int_0^y d\nu \bar{V}(\bar{x},\bar{y};x,y) e^{-(C_2\bar{x}+C_1\bar{y})} \\ \bar{E}(x,y)_{\mathbf{x}} &= K_5 \left[\int_0^y d\nu \bar{V}(x,\bar{y};x,y) e^{-(C_2x+C_1\bar{y})} \\ &+ \int_0^x d\mu \int_0^y d\nu \bar{V}(\bar{x},\bar{y};x,y)_{\mathbf{x}} e^{-(C_2\bar{x}+C_1\bar{y})} \right] \end{split}$$

(3.22)

$$\begin{split} \bar{E}(x, y)_{Y} &= K_{5} \left[\int_{0}^{x} d\mu \bar{V}(\bar{x}, y; x, y) e^{-(C_{2}\bar{x}+C_{1}\bar{y})} \\ &+ \int_{0}^{x} d\mu \int_{0}^{y} d\nu \bar{V}(\bar{x}, \bar{y}; c, y)_{Y} e^{-(C_{2}\bar{x}+C_{1}\bar{y})} \right] \\ \bar{E}(x, y)_{XY} &= K_{5} e_{h}^{-(C_{2}x+C_{1}y)} + C_{5} \bar{E}(x, y). \end{split}$$

$$E(x, y)_{xy} = K_5 e_h + C_5 E(x, y)_{xy}$$

From (3.21) and (3.22) it follows that

$$\begin{split} \bar{E}(x,y) &\leq K_5 \int_0^x d\mu \int_0^y d\nu e_h^{\sqrt{C_5}(x-\bar{x}+y-\bar{y})} e_h^{-(C_2\bar{x}+C_1\bar{y})} = \beta_1 \\ \bar{E}_x(x,y) &\leq K_5 \int_0^y d\nu e_h^{-[C_2x+C_1\bar{y}]} + \sqrt{C_5} \int_0^x d\mu \int_0^y d\nu e_h^{\sqrt{C_5}(x-\bar{x}+y-\bar{y})} e_h^{-(C_2\bar{x}+C_1\bar{y})} = \beta_2 \\ \bar{E}_x(x,y) &\leq K_5 \int_0^x d\mu e_h^{-[C_2\bar{x}+C_1\bar{y}]} + \sqrt{C_5} \int_0^x d\mu \int_0^y d\nu e_h^{\sqrt{C_5}(x-\bar{x}+y-\bar{y})} e_h^{-(C_2\bar{x}+C_1\bar{y})} = \beta_3 \\ \bar{E}(x,y)_{xy} &\leq K_5 e_h^{-[C_2x+C_1\bar{y}]} + C_5 \beta_1 . \\ \text{Thus from (3.14) and (3.16) we obtain} \\ &|E(x,y)| \leq E^*(x,y) \leq \beta_1 e^{C_2x+C_1\bar{y}} = \gamma_1 \\ &|E(x,y)| \leq E^*(x,y) \leq \kappa^{C_2x+C_1\bar{y}} = \kappa^{C_2x+C_1\bar{y}} = \gamma_1 \end{split}$$

$$|E(x, y)_{\mathbf{x}}| \leq E^{*}(x, y)_{\mathbf{x}} \leq e_{h}^{2 - r - r} [\beta_{2} + C_{2}\beta_{1}] = \gamma_{2}$$

$$(3.23) |E(x, y)_{\mathbf{x}}| \leq E^{*}(x, y)_{\mathbf{x}} \leq e_{h}^{c_{2}x + c_{1}y} [\beta_{3} + C_{1}\beta_{1}] = \gamma_{3}$$

$$|E(x, y)_{\mathbf{x}\mathbf{x}}| \leq E^{*}(x, y)_{\mathbf{x}\mathbf{x}} \leq K_{5} + e_{h}^{c_{2}x + c_{1}y} [C_{5}\beta_{1} + C_{1}\beta_{2} + C_{1}C_{2}\beta_{1} + C_{2}\beta_{3}]$$

$$= \gamma_{4}$$

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Now substitution of the bounds in (3.23) in (3.12) yields

(3.24)
$$|u_x(P) - U(P)_x| \leq \gamma_2 + h^2 \alpha_2, \qquad |u_y(P) - U(P)_Y| \leq \gamma_3 + h^2 \alpha_2, \\ |u_{xy}(P) - U(P)_{XY}| \leq \gamma_4 + h^2 \alpha_1.$$

4. General Quasi-Linear Case. In this section we shall consider a method for approximating the solution of

(4.1)
$$u_{xy} = f(x, y, u, v, w), \qquad (x, y) \in R$$
$$v = u_x, \qquad w = u_y$$
$$u(x, 0) = u(0, y) = 0.$$

We assume that f(x, y, u, v, w) satisfies the Lipschitz condition

(4.2)
$$|f(x, y, \bar{u}, \bar{v}, \bar{w}) - f(x, y, u, v, w)| \\ \leq K_1 |\bar{u} - u| + K_2 |\bar{v} - v| + K_3 |\bar{w} - w|,$$

and is bounded for

$$-\infty < u, v, w < \infty$$
.

Under these conditions it is known [Diaz, 6] that (4.1) possesses a unique solution.

We introduce a network as before and define the difference quotients as in (2.4). Again we extend the domain of definition of an arbitrary mesh function to the centers of squares as averages so that the chain rule for differencing (2.6) is valid. We define the finite difference analogue of (4.1) to be

(4.3)

$$U_{XY}(P) = f(P, U(P), U_X(P), U_Y(P)),$$

$$P = (mh + \frac{1}{2}h, nh + \frac{1}{2}h) \qquad m, n = 0, 1, 2, 3, \cdots$$

$$U(x, 0) = U(0, y) = 0.$$

At each square with P given in (4.3) the difference equation becomes

(4.4)
$$U[(m+1)h, (n+1)h] = U(mh, nh) + h[U_x(mh + \frac{1}{2}h, nh) + U_y(mh, nh + \frac{1}{2}h)] + h^2 f[P, U(P), U_x(P), U_y(P)].$$

Then (4.4) is of the form

$$(4.5) U(mh, nh) = g[U(mh, nh)]$$

where from (4.2) g is a continuous function of U(mh, nh). If we assume that h is so small that

(4.6)
$$\frac{1}{2}h[K_2 + K_3 + \frac{1}{2}hK_1] < 1$$

then by the contraction mapping theorem we can solve (4.4) for U[(m + 1)h, (n + 1)h] in terms of U[(m + 1)h, nh], U[mh, (n + 1)h], and U(mh, nh), through an iterative process.

The discretization error E(P) = u(P) - U(P) satisfies the difference inequality

(4.7)
$$E_{XY} \leq K_{1} |E| + K_{2} |E_{X}| + K_{3} |E_{Y}| + K_{4}$$
$$K_{4} = K_{1} |u - \bar{u}|_{M} + K_{2} |u_{x} - u_{X}|_{M}$$
$$+ K_{3} |u_{y} - u_{Y}|_{M} + |u_{xy} - u_{XY}|_{M}$$

at points $P = (mh + \frac{1}{2}h, nh + \frac{1}{2}h)$ with

$$\bar{u}(P) = \frac{1}{4} \{ u[(m+1)h, (n+1)h] + u[(m+1)h, nh] \}$$

$$+ u[mh, (n + 1)h] + u(mh, nh)$$

In addition we note that

$$E(x, 0) = E(0, y) = 0.$$

We pose the majorizing problem

(4.8)
$$W_{xy} = K_1 W + K_2 W_x + K_3 W_y + K_4 W_y + K$$

By an obvious modification of the proof of Theorem 2.1, Corollary 2 we conclude that

$$(4.9) |E| \leq W, |E_{\mathbf{X}}| \leq W_{\mathbf{X}}, |E_{\mathbf{Y}}| \leq W_{\mathbf{Y}}, |E_{\mathbf{X}\mathbf{Y}}| \leq W_{\mathbf{X}\mathbf{Y}}.$$

The following theorem gives the bound on the discretization error in this case.

THEOREM 4.1. Let E = u - U where u and U are the solutions of (4.1) and (4.3) respectively. Then E has the bound

(4.10)
$$|E| \leq K_4 \int_0^{\xi} dx \int_0^{\eta} dy [e_h^{\sqrt{C_5}\{(\xi-x)+(\eta-y)\}}] e_h^{K_3(\xi-x)+K_2(\eta-y)}$$

where C_5 is given by

$$(4.11) C_5 = K_1 + K_2 K_3$$

Proof: The theorem follows from (4.9) and Theorem 2.1, Corollary 1.

The bound (4.10) on the discretization error will be explicit once we show how to bound K_4 in terms of the data of the problem. At this point we assume that the partial derivatives of f with respect to u, v, and w through the third order are bounded. The quantity K_4 in (4.7) can be bounded in terms of the higher derivatives of u(x, y) as is seen from (3.4). Now if $|f|_M = |f(x, y, 0, 0, 0)|_M$ and v(x, y)is a solution of

(4.12)
$$v_{xy} = K_1 v + K_2 v_x + K_3 v_y + |f|_M$$
$$v(x, 0) = v(0, y) = 0$$

then by the method of successive approximations we see that

$$(4.13) |u| \leq v, |u_x| \leq v_x, |u_y| \leq v_y, |u_{xy}| \leq v_{xy},$$

Hence we see that bounds for u, u_x, u_y, u_{xy} , are given by (3.7) where we make the substitutions

$$(4.14) |a|_{M} = K_{2}, |b|_{M} = K_{3}, \beta = K_{1} + K_{2}K_{3}.$$

Bounds for the higher mixed derivatives in terms of the known bounds for |u|, $|u_x|$, $|u_y|$, $|u_{xy}|$ are obtained by differentiating the differential equation (4.1), as was done in section 3. To obtain bounds for the higher pure derivatives we use a variation of the method applied in section 3. We illustrate the method for u_{xx} .

From (4.1) we see that

$$(4.15) u_{xxy} = f_x + f_u u_x + f_v u_{xx} + f_w u_{xy}$$

Hence by introducing an integrating factor

(4.16)
$$u_{xx}(x,y) = \int_0^y d\eta g(x,\eta) \exp\left[-\int_\eta^y l(x,\lambda) d\lambda\right]$$

where

(4.17)
$$g(x, y) = f_x + f_u u_x + f_w u_{xy}$$
$$l(x, y) = f_v[x, y, u(x, y), u_x(x, y), u_y(x, y)].$$

Now under the assumption that f_x , f_u , f_v , f_w are bounded we can obtain an explicit bound for u_{xx} . The bounds thus obtained for the higher derivatives allow us to bound K_i in (4.10).

5. An example. Consider the boundary value problem:

(5.1)
$$u_{xy}(x, y) - xyu(x, y) = x^2y^2(9 - x^2y^2)$$
$$u(x, 0) = u(0, y) = 0$$

in the closed square $s: \{0 \le x \le 1, 0 \le y \le 1\}$, which has the solution $u = x^3 y^3$. From (3.7) we have

$$| u(\xi, \eta) |_{M} \leq \int_{0}^{\xi} \int_{0}^{\eta} | v(x, y; \xi, \eta) | (9 - x^{2}y^{2})x^{2}y^{2} dy dx$$

$$| u(\xi, \eta)_{\xi} |_{M} \leq \int_{0}^{\eta} | v(\xi, y; \xi, \eta) | (9 - \xi^{2}y^{2}) dy dx$$

(5.2)

$$+ \int_{0}^{\xi} \int_{0}^{\eta} | v_{\xi}(x, y; \xi, \eta) | (9 - x^{2}y^{2})x^{2}y^{2} dy dx$$

$$| u(\xi, \eta)_{\xi} |_{M} \leq \int_{0}^{\xi} | v(x, \eta; \xi, \eta) | (9 - x^{2}\eta^{2})x^{2}\xi^{2} dx$$

$$+ \int_{0}^{\xi} \int_{0}^{\eta} | v_{\eta}(x, y; \xi, \eta) | (9 - x^{2}y^{2})x^{2}y^{2} dy dx$$

where

$$\begin{aligned} v(x, y; \xi, \eta) &= I_0[\sqrt{|xy|_{\eta}(\xi - x)(\eta - y)}] \leq I_0(2) \approx 2.28\\ v(\xi, y; \xi, \eta) &= v(x, \eta; \xi, y) = 1 \end{aligned}$$

(5.3)
$$v_{\xi}(x,y;\xi,\eta) = \frac{\sqrt{\eta-y}}{\sqrt{\xi-x}} I_1[2\sqrt{|xy|_{\eta}(\xi-x)(\eta-y)}]$$
$$\leq \frac{\sqrt{\eta-y}}{\sqrt{\xi-x}} I_1(2) \approx 1.6 \frac{\sqrt{\eta-y}}{\sqrt{\xi-x}}$$

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(5.4)
$$v_{\eta}(x,y;\xi,\eta) \leq \frac{\sqrt{\xi-x}}{\sqrt{\eta-y}} I_1(2) \approx 1.6 \frac{\sqrt{\xi-x}}{\sqrt{\eta-y}}$$

From (5.2) and (5.3) by integration it follows that

$$|u| \leq 2.18; |u_x|, |u_y| \leq 5.36$$

$$|u_{xy}| \leq 11.18; |u_{xx}|, |u_{yy}| \leq 25.54$$

$$(5.5) |u_{xxy}|, |u_{xyy}| \leq 25.54; |u_{xxx}|, |u_{yyy}| \leq 54.26$$

$$|u_{xxxy}|, |u_{xyyy}| \leq 54.26; |u_{xxyy}| \leq 59.08$$

$$|u_{xxxx}|, |u_{yyyy}| \leq 148.88$$

Now $K_4(1, 1)$ appearing in (3.2) is given by

(5.6)
$$K_4(1, 1) \leq h^2(\alpha_1 + \alpha_3)$$

where α_1 and α_3 are given by (3.4). We have $\alpha_1 = 11.36$, $\alpha_3 = 9.18$, thus from (3.2) we obtain

(5.7)
$$|E|_{M} \leq 20.54h^{2} \int_{0}^{\xi} \int_{0}^{\xi} e_{\eta} \sqrt{|xy|_{\eta}\{(\xi - x) + (\eta - y)\}} dx dy \leq 60.59h^{2}$$

The solution to the finite difference analog of (5.1) was obtained on the computer at the U. S. Naval Ordnance Laboratory. The calculations were made with h = 0.1 with the results given in the following table.

		· · · · · · · · · · · · · · · · · · ·		1
x	y	u(x, y)	U(x, y)	E = u - U
.2	.2	$6.40 imes10^{-5}$	6.392×10^{-5}	$8.0 imes 10^{-8}$
.2	.4	5.12×10^{-4}	5.116×10^{-4}	4.0×10^{-7}
.2	.6	$1.728 imes 10^{-3}$	$1.727 imes 10^{-3}$	$1.2 imes 10^{-6}$
.2	.8	$4.096 imes 10^{-3}$	4.093×10^{-3}	$2.7 imes 10^{-6}$
.2	1.0	$8.0 imes 10^{-3}$	7.995×10^{-3}	$5.2 imes10^{-6}$
.4	.4	4.096×10^{-3}	4.095×10^{-3}	1.3×10^{-6}
.4	.6	1.3824×10^{-2}	1.3821×10^{-2}	3.2×10^{-6}
.4	.8	3.2768×10^{-2}	3.2762×10^{-2}	$6.5 imes 10^{-6}$
.4	1.0	$6.4 imes 10^{-2}$	6.3988×10^{-2}	1.2×10^{-5}
.6	.6	4.6656×10^{-2}	4.6649×10^{-2}	6.7×10^{-6}
.6	.8	1.1059×10^{-1}	1.1059×10^{-1}	1.3×10^{-5}
.6	1.0	2.16×10^{-1}	2.1598×10^{-1}	$2.2 imes 10^{-5}$
.8	.8	2.6214×10^{-1}	2.6212×10^{-1}	2.2×10^{-5}
.8	1.0	5.12×10^{-1}	5.1196×10^{-1}	3.7×10^{-5}
1.0	1.0	1.0	9.9994×10^{-1}	6.2×10^{-5}

In this case the *a priori* bound on the discretization error given by 5.7 is 6×10^{-3} and the observed maximum error in the table above is 6×10^{-5} . The fact that the error bound is 100 times too large is in part attributable to the crude manner in which the higher derivatives of u(x, y) have been computed. It is clear that the higher derivatives should be solved for directly in terms of u, u_x, u_y for the particular problem under consideration with inequalities obtained only at the last moment.

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